

Remarks on statistical errors in equivalent widths*

K. Vollmann and T. Eversberg**

Schnörringen Telescope Science Institute, Ringweg 7, 51545 Waldbröl, Germany

Received 2006 Mar 29, accepted 2006 Jun 14

Published online 2006 Oct 16

Key words line: profiles – methods: statistical – techniques: spectroscopic

Equivalent width measurements for rapid line variability in atomic spectral lines are degraded by increasing error bars with shorter exposure times. We derive an expression for the error of the line equivalent width $\sigma(W_\lambda)$ with respect to pure photon noise statistics and provide a correction value for previous calculations.

© 2006 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim

1 Introduction

Error bars of line equivalent widths W_λ have been the subject of a number of investigations of high temporal and spectral resolution observations with exposures of the order of minutes or even seconds, leading to large errors. A first expression for the calculation of $\sigma(W_\lambda)$ for pure photon noise statistics was estimated by Williams et al. (1974) for their rapid line variability observations of Ap and Bp stars. This was then supplemented by Young (1974) for errors due to scintillation. Lacy (1977) published an improvement of Williams et al. and compared pure photon noise statistics with the case including scintillation for his H α and H β rapid line investigations of massive stars. Guided by his work, Chalabaev & Maillard (1983) again derived an expression for the error from pure photon noise statistics in an appendix of their paper on rapid spectral variability of Be stars. However, their derivation of $\sigma(W_\lambda)$ shows some inconsistencies which we attempt to resolve in this paper.

2 The equivalent width of a spectroscopic line

In practice, the measurement of noise within a line can be a difficult task, because photon noise is superposed with stellar variations. Hence, it is necessary to estimate an expression for the equivalent width which separates these two parameters and then to find an expression for its error $\sigma(W_\lambda)$. We start with the standard definition of the equivalent width:

$$W_\lambda = \int_{\lambda_1}^{\lambda_2} \frac{F_c(\lambda) - F(\lambda)}{F_c(\lambda)} d\lambda$$

or

$$W_\lambda = \int_{\lambda_1}^{\lambda_2} \left[1 - \frac{F(\lambda)}{F_c(\lambda)} \right] d\lambda, \quad (1)$$

* This publication is dedicated to Karl-Werner Eversberg who passed away in October 2005.

** Corresponding author: eversberg@stsci.edu

with $F_c(\lambda)$ the flux in the continuum, $F(\lambda)$ the flux in the line at the wavelength λ and $F(\lambda) = F_c(\lambda)$ for $\lambda \geq \lambda_2$ and $\lambda \leq \lambda_1$. In a first step we integrate Eq. (1) separately and we obtain

$$W_\lambda = \Delta\lambda - \int_{\lambda_1}^{\lambda_2} \frac{F(\lambda)}{F_c(\lambda)} d\lambda,$$

with the wavelength interval $\Delta\lambda = \lambda_2 - \lambda_1$. By applying the mean value theorem we obtain

$$W_\lambda = \Delta\lambda \cdot \left[1 - \overline{\left(\frac{F(\lambda)}{F_c(\lambda)} \right)} \right], \quad (2)$$

with $\overline{\left(\frac{F(\lambda)}{F_c(\lambda)} \right)}$ the normalized average flux within $\Delta\lambda$. Equation (2) can be written as

$$W_\lambda \approx \Delta\lambda \cdot \left[1 - \frac{\overline{F}}{F_c} \right]. \quad (3)$$

This can be shown by applying a discrete approximation of Eq. (2). We use the arithmetic mean and substitute the line and continuum fluxes by their average values plus their deviations ΔF_i and $\Delta F_{c,i}$:

$$\overline{\left(\frac{F}{F_c} \right)} = \frac{1}{M} \cdot \sum_{i=1}^M \frac{F_i}{F_{c,i}} \Leftrightarrow \overline{\left(\frac{F}{F_c} \right)} = \frac{1}{M} \cdot \sum_{i=1}^M \left(\frac{\overline{F} + \Delta F_i}{\overline{F_c} + \Delta F_{c,i}} \right)$$

M represents the number of pixels within the line. Note that ΔF_i within the line consists not only of noise but also of line information and can be quite large in contrast to $\Delta F_{c,i}$,

$$\overline{\left(\frac{F}{F_c} \right)} = \frac{\overline{F}}{\overline{F_c}} \cdot \frac{1}{M} \cdot \sum_{i=1}^M \frac{(1 + \frac{\Delta F_i}{\overline{F}})}{(1 + \frac{\Delta F_{c,i}}{\overline{F_c}})}.$$

In case of sufficient S/N within the continuum we have

$$\left| \frac{\Delta F_{c,i}}{\overline{F_c}} \right| \ll 1 \Rightarrow 1 + \frac{\Delta F_{c,i}}{\overline{F_c}} \approx 1,$$

$$\overline{\left(\frac{F}{F_c} \right)} = \frac{\overline{F}}{\overline{F_c}} \cdot \frac{1}{M} \cdot \sum_{i=1}^M \left(1 + \frac{\Delta F_i}{\overline{F}} \right),$$

$$\overline{\left(\frac{F}{F_c}\right)} = \frac{\overline{F}}{\overline{F_c}} + \frac{1}{\overline{F_c}} \cdot \frac{1}{M} \cdot \sum_{i=1}^M \Delta F_i.$$

However, by definition

$$\frac{1}{M} \cdot \sum_{i=1}^M \Delta F_i = 0 \implies \overline{\left(\frac{F}{F_c}\right)} \approx \frac{\overline{F}}{\overline{F_c}}.$$

Hence, the equivalent width can be estimated from the wavelength interval and the average intensities in the line and the continuum.

3 The error of the equivalent width

The intensities $F(\lambda)$ and $F_c(\lambda)$ are influenced by certain statistical errors. If we keep in mind that F_c is generally estimated outside the line region and interpolated across the line where the line flux is measured and if we assume that the errors are not correlated we are able to estimate their standard deviations separately. Following the principal of error propagation we expand Eq. (3) in a Taylor series:

$$W_\lambda = W(\overline{F}, \overline{F_c}) + \frac{\partial W_\lambda}{\partial \overline{F}} (F - \overline{F}) + \frac{\partial W_\lambda}{\partial \overline{F_c}} (F_c - \overline{F_c}),$$

where F and F_c are random variables.

The variance $Var(W_\lambda) \equiv \sigma^2(W_\lambda)$ of the expansion is

$$\sigma^2(W_\lambda) = \left[\frac{\partial W_\lambda}{\partial \overline{F}} \cdot \sigma(F) \right]^2 + \left[\frac{\partial W_\lambda}{\partial \overline{F_c}} \cdot \sigma(F_c) \right]^2, \quad (4)$$

with $\sigma(F)$ and $\sigma(F_c)$ the standard deviations in the line and the continuum, respectively. By using Eq. (3) the two partial derivatives are

$$\frac{\partial W_\lambda}{\partial \overline{F}} = -\frac{\Delta \lambda}{\overline{F_c}},$$

with $\Delta \lambda = M \cdot h_\lambda$ and the spectral dispersion h_λ , and

$$\frac{\partial W_\lambda}{\partial \overline{F_c}} = \frac{1}{\overline{F_c}} (\Delta \lambda - W_\lambda).$$

3.1 Weak lines

For weak lines the depth of the line is negligible and

$$\sigma(F) = \frac{\overline{F}}{S/N} \cong \sigma(F_c) = \frac{\overline{F_c}}{S/N}$$

with the signal-to-noise ratio S/N in the undisturbed continuum. With these terms we can write

$$\sigma^2(W_\lambda) = \left[\frac{M \cdot h_\lambda}{S/N} \right]^2 \cdot \left[\frac{\overline{F}}{\overline{F_c}} \right]^2 + \left[\frac{\sigma(F_c)}{\overline{F_c}} \cdot (\Delta \lambda - W_\lambda) \right]^2, \quad (5)$$

where the first term is the photometric uncertainty and the second term the uncertainty of the continuum estimation over the line. According to Eq. (3) both terms are identical and we have

$$\sigma(W_\lambda) = \sqrt{2} \cdot \frac{(\Delta \lambda - W_\lambda)}{S/N}. \quad (6)$$

This result is easy to understand if we realize that the line strength for weak lines will be of the order of the noise itself.

3.2 Low- and high-flux lines

In the case of low (absorption) and high (emission) flux lines we have to use the corresponding Poisson-statistic

$$\sigma(F) = \sqrt{\frac{\overline{F}}{\overline{F_c}}} \cdot \sigma(F_c),$$

and we get

$$\sigma(W_\lambda) = \sqrt{1 + \frac{\overline{F_c}}{\overline{F}}} \cdot \frac{(\Delta \lambda - W_\lambda)}{S/N}. \quad (7)$$

This result seems to be quite simple. However, we should not expect a complex solution for the error bars of equivalent widths because of a relatively simple definition of W . In both cases we now can obtain the standard deviations with the measurable parameters S/N , $\Delta \lambda$, \overline{F} and $\overline{F_c}$. In addition, Eq. (7) represents the general error of the line equivalent width and for the case $\overline{F} \approx \overline{F_c}$ we obtain the result for weak lines again.

4 A comparison with former calculations

Interestingly our result in Eq. (5) differs from Eq. (A9) of Chalabaev & Maillard (1983, ‘‘C&M’’) by a factor M , the number of pixels corresponding to the interval $\lambda_1 \leq \lambda \leq \lambda_2$, in the first summand. The reason for this difference can be found in their expansion C&M (A3):

- C&M (A3) is a multidimensional expansion for the equivalent width with average values $\overline{F_j}$ for all pixel fluxes F_j (the pixel fluxes F_j define random variables). However, it is not clear why the expansion over M fluxes F_j should be necessary.
- Although there is no principal difference between the line and the continuum the continuum part is not consequently developed as it is done for the line part. They also say that ‘‘the value (S/N) characterizes the spread of read-outs on the continuum of an individual spectrum, while the quantity $\sigma(\overline{F_c})$ is the total uncertainty of the continuum level determination, and these two quantities may be related in a non-simple way’’. However, we assume that the uncertainty of the continuum is simply defined by the noise within the continuum and therefore the above distinction is not necessary.
- Another difference between C&M and our approach can be found in their Eq. (A7) which gives $S/N = \frac{F_j}{\sigma(F_j)}$. However, the correct expression is $S/N = \frac{\overline{F}}{\sigma(F)}$ (which is $\frac{\overline{F_j}}{\sigma(F_j)}$ in the C&M notation).

4.1 The \sqrt{N} -rule

Generally, for the variance of a random variable x we have $Var(x) = \overline{x^2} - \overline{x}^2$ with \overline{x} the expectancy value of x . Respectively, we get $Var(k \cdot x) = k^2 \cdot Var(x)$ with $k = constant$.

In particular, the mean value of \bar{x} of N random variables x_1, \dots, x_N is $\bar{x} = \frac{1}{N} \cdot \sum_{i=1}^N x_i$ and hence

$$\text{Var}(\bar{x}) = \text{Var}\left(\frac{1}{N} \cdot \sum_{i=1}^N x_i\right) = \frac{1}{N^2} \cdot \sum_{i=1}^N \text{Var}(x_i).$$

Note that \bar{x} represents a random variable, as well. If all variances $\text{Var}(x_i)$ are identical the \sqrt{N} -rule gives $\sigma(\bar{x}) = \sigma(x)/\sqrt{N}$.

In their approach C&M develop W_λ in a multi-dimensional expansion

$$W_\lambda = \bar{W} + \sum_{j=1}^M \frac{\partial W_\lambda}{\partial F_j} \cdot (\bar{F}_j - F_j) + \frac{\partial W_\lambda}{\partial F_c} \cdot (\bar{F}_c - F_c).$$

We interpret the C&M notation of \bar{F}_j and \bar{F}_c as follows: $\bar{F}_j = \frac{1}{M} \cdot \sum_{j=1}^M F_j$ and $\bar{F}_c = \frac{1}{\tilde{N}} \cdot \sum_{i=1}^{\tilde{N}} F_{c,i}$ with \tilde{N} the number of pixel fluxes for the determination of the noise. Using $\frac{\partial W_\lambda}{\partial F_j} = -\frac{M \cdot h_\lambda}{F_c \cdot M}$ they get

$$W_\lambda = -\frac{M \cdot h_\lambda}{F_c} \cdot \frac{1}{M} \cdot \sum_{j=1}^M (\bar{F}_j - F_j) + \frac{\partial W_\lambda}{\partial F_c} \cdot (\bar{F}_c - F_c).$$

By calculating the variance of W_λ one obtains

$$\begin{aligned} \text{Var}(W_\lambda) &= \left(\frac{M \cdot h_\lambda}{F_c}\right)^2 \cdot \text{Var}\left(\frac{1}{M} \cdot \sum_{j=1}^M (\bar{F}_j - F_j)\right) \\ &+ \left(\frac{\partial W_\lambda}{\partial F_c}\right)^2 \cdot \text{Var}(\bar{F}_c - F_c). \end{aligned}$$

In their approach C&M now assume that both average values \bar{F} and \bar{F}_c are constant (although this assumption is just an approximation for \bar{F} by definition) and hence one gets

$$\begin{aligned} \text{Var}(W_\lambda) &= \left(\frac{M \cdot h_\lambda}{F_c}\right)^2 \cdot \text{Var}\left(\frac{1}{M} \cdot \sum_{j=1}^M F_j\right) \\ &+ \left(\frac{\partial W_\lambda}{\partial F_c}\right)^2 \cdot \text{Var}(F_c). \end{aligned}$$

Now it is clear that C&M's goal was to estimate the variance within the first summand with the variance over the averages $\bar{F} = \frac{1}{M} \cdot \sum_{j=1}^M F_j$. The \sqrt{N} -rule gives

$$\begin{aligned} \text{Var}(W_\lambda) &= \left(\frac{M \cdot h_\lambda}{F_c}\right)^2 \cdot \frac{1}{M} \cdot \text{Var}(F_j) \\ &+ \left(\frac{\Delta\lambda - W_\lambda}{F_c}\right)^2 \cdot \text{Var}(F_c), \end{aligned}$$

and hence for the weak-line case

$$\begin{aligned} \text{Var}(W_\lambda) &= \left(\frac{M \cdot h_\lambda}{S/N}\right)^2 \cdot \left(\frac{\bar{F}}{F_c}\right)^2 \cdot \frac{1}{M} \\ &+ \frac{\text{Var}(F_c)}{F_c^2} \cdot (\Delta\lambda - W_\lambda)^2. \end{aligned}$$

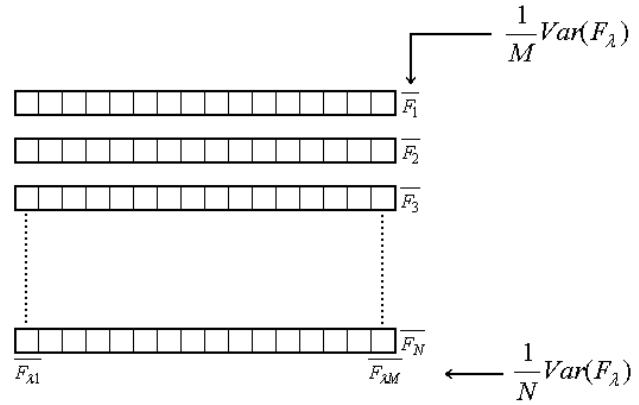


Fig. 1 Principle of the variance and average determination for a single spectrum and a number of spectra. The figure shows N measured spectra with M pixels each.

This is C&M (A9), and the first summand is smaller than our result in Eq. (5) by the factor M . For clarification we show the principal procedures in Fig. 1. Every measured single spectrum delivers a separate variance $\text{Var}(F_\lambda)$. By obtaining N spectra, the variance for the wavelength-averaged spectrum is $\text{Var}(\bar{F}_\lambda) = \frac{1}{N} \cdot \text{Var}(F_\lambda)$. But considering the variance of all spectral averages, as is done by C&M, we get $\text{Var}(\bar{F}) = \frac{1}{M} \cdot \text{Var}(F_\lambda)$ independent of the number of measured spectra. Because of the missing sum within the continuum part of C&M (A3) the factor $1/M$ does not appear in both sums.

If we want to estimate $\sigma(W_\lambda)$ properly by means of C&M (A9), depending on the S/N over the whole line, we have to scale up the measured noise by the factor \sqrt{M} . Therefore the corrected formula C&M (A7) is $S/N = \sqrt{M} \cdot \frac{F_j}{\sigma(F_j)}$. By using this correction Chalabaev & Maillard (1983) would have found the same result for $\sigma(W_\lambda)$ as we do. However, we note that this correction is enforced by the series expansion of C&M. Their procedure, however, has been used in a series of papers (more than 40 until now) using expression (A9) of Chalabaev & Maillard (1983).

Now it is easy to estimate the correction factor for the hitherto estimated error bars of W_λ . For the normal case that both parts $\sigma(F)$ and $\sigma(F_c)$ are identical we find

$$\sigma(W_\lambda) = \sqrt{\frac{2M}{M+1}} \cdot \sigma(W_\lambda)_{\text{C\&M}}.$$

For broad lines we have $M+1 \approx M$ and hence $\sigma(W_\lambda) \approx \sqrt{2} \cdot \sigma(W_\lambda)_{\text{C\&M}}$ where the subindex C&M indicates the values obtained by Chalabaev & Maillard (1983). In addition, if we assume that both errors (line and continuum) are different (as done by C&M) our result can differ even more.

In fact, the above Eq. (4) for $\sigma^2(W_\lambda)$ does not represent the exact variance of W but is just an approximation, which is especially important in the case of small basic populations (i.e. the number of pixel fluxes for the determination of noise). The exact definition of the unbiased variance of

W_λ is

$$\sigma^2(W_\lambda) = \frac{\tilde{N} \cdot (M - 1)}{\tilde{N} \cdot M - 1} \cdot \left(\frac{\partial W_\lambda}{\partial F} \cdot \sigma(F) \right)^2 + \frac{M \cdot (\tilde{N} - 1)}{\tilde{N} \cdot M - 1} \cdot \left(\frac{\partial W}{\partial F_c} \cdot \sigma(F_c) \right)^2 \quad (8)$$

(Kreyszig 1967) with \tilde{N} the number of pixel values for the estimation of S/N and M the number of pixels within the line. This equation transforms into Eq. (4) if \tilde{N} and M become sufficiently large. If the number of pixels for the estimation of S/N and/or the line width are too small, the error of equivalent width will be overestimated by using Eq. (4). This is especially important for low-resolution observations. For example, if the investigated line covers only 10 pixel and S/N is estimated over about 50–100 pixels, the true error would be overestimated by about 3%.

4.2 Scrutinizing previous results

Former conclusions regarding W_λ have to be carefully scrutinized due to increased error bars with respect to the result of C&M, especially if the interpretation of the data depends on results at the detection limit. For clarification, we give two prominent examples:

- Searching for rapid line variability in the Be star BD +37°1160, Rossi et al. (1991) tried to find a W_λ correlation between H α and H β (see their Fig. 4). Their conclusion of “...spectral line variability on a time scale of 300 s” for this Be star is unsustainable, especially for H β .
- Investigating periodic variabilities in UV lines of the B star ζ Cas, Neiner et al. (2003) claimed to see minima and maxima for W_λ of the NV 1240 doublet (see their Fig. 2). This is not valid with respect to our correction factor of $\sqrt{2}$.

Acknowledgements. The authors would like to thank Tony Moffat and Günter Gebhard for helpful discussions. K.V. would like to thank his wife Susanne for her patience. T.E would like to thank Karin and Karl-Werner Eversberg for their support.

References

- Chalabaev, A., Maillard, J.P.: 1983, A&A 127, 279
 Kreyszig, E.: 1967, *Statistische Methoden und ihre Anwendungen*, Göttingen; Vandenhoeck & Ruprecht, p.333
 Lacy, D.: 1977, ApJ 212, 132
 Neiner, C., Geers, V.C., Henrichs, H.F., Floquet, M., Frémat, Y., Hubert, A.-M., Preuss, O., Wiersema, K.: 2003, A&A 406, 1019
 Rossi, C., Norci, L., Polcaro, V.F.: 1991, ApJL 249, L19
 Williams, W.E., Frantz, R.L., Breger, M.: 1974, A&A 35, 381
 Young, A.G.: 1974, in: N. Carleton (ed.), *Methods of Experimental Physics: Astrophysics*, New York, Academic Press, p. 101